

From symplectic to spin geometry

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Quantization for spinless system

| classical | quantum | link |
|------------------------------------------------------------------------------|---------------------------------------------------------------------------------------------------------|--------------------------------------------------------------------------------------------------------------------------------------------------------------------------|
| symplectic mfld (T^*M, ω) | (pre-)Hilbert space $\mathcal{F}_c^{\frac{1}{2}} = \Gamma_c(\Lambda^n T^*M ^{\otimes \frac{1}{2}})$ | vertical polarization |
| co. graded P. alg. $\text{Pol}(T^*M)$ | ass. filtered alg. $\mathcal{D}^{\frac{1}{2}, \frac{1}{2}}(M)$ | $\text{Pol}_k \simeq \mathcal{D}_k^{\frac{1}{2}, \frac{1}{2}} / \mathcal{D}_{k-1}^{\frac{1}{2}, \frac{1}{2}}$ $\sigma_{p+q-1}([P, Q]) = \{\sigma_p(P), \sigma_q(Q)\}$ |
| preserve $\{\cdot, \cdot\}$ $\text{Vect}(M) \hookrightarrow \text{Pol}_1$ | preserve $[\cdot, \cdot]$ $\text{Vect}(M) \hookrightarrow \mathcal{D}_1^{\frac{1}{2}, \frac{1}{2}}$ | $\text{GQ}(J_X) = X^i \partial_i + \frac{1}{2}(\partial_i X^i).$ |

Quantization for spin system

Let (M, g) be a pseudo-Riemannian spin manifold with $\dim M = 2n$.
Remind that $\mathbb{Cl}(V, g) = \bigotimes V / \langle u \otimes v + v \otimes u + 2g(u, v) \rangle$.

$$\begin{array}{c|c} & \Gamma(S) \otimes \mathcal{F}_c^{\frac{1}{2}} \\ \hline \text{Pol}(T^*M) \otimes \Gamma(\mathbb{Cl}(M, g)) & D^{\frac{1}{2}, \frac{1}{2}} \ni P_{j_1 \dots j_k}^{i_1 \dots i_k}(x) \gamma^{j_1} \dots \gamma^{j_k} \partial_{i_1} \dots \partial_{i_k} \end{array}$$

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$$\begin{array}{c|c} & \Gamma(S) \otimes \mathcal{F}_c^{\frac{1}{2}} \\ \hline Sb=\text{Pol}(T^*M) \otimes \Omega_{\mathbb{C}}(M) & D^{\frac{1}{2}, \frac{1}{2}} \ni P_{j_1 \dots j_k}^{i_1 \dots i_k}(x) \gamma^{j_1} \dots \gamma^{j_k} \partial_{i_1} \dots \partial_{i_k} \end{array}$$

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$$\mathcal{M} = T^*M \oplus \Pi TM$$

$$\Gamma(S) \otimes \mathcal{F}_c^{\frac{1}{2}}$$

$$\text{Sb}=\text{Pol}(T^*M) \otimes \Omega_{\mathbb{C}}(M)$$

$$D^{\frac{1}{2}, \frac{1}{2}} \ni P_{j_1 \dots j_\kappa}^{i_1 \dots i_k}(x) \gamma^{j_1} \dots \gamma^{j_\kappa} \partial_{i_1} \dots \partial_{i_k}$$

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Questions :

- symplectic form on \mathcal{M} ?
- Polarization and construction of S ?
- Quantum action of vector fields ?
- Poisson bracket from the commutator, which graduation ?

1 From symplectic spinors to spinors

- Algebraic point of view
- Geometrization

2 Actions of vector fields

- On the supercotangent bundle \mathcal{M}
- On the spinor bundle S

3 D as quantization of Sb

Quantization of symplectic even and odd vector spaces

Let V be a vector space, with coordinates (ξ^i) , $\dim V = 2n$. Remind that $\mathcal{W}(V, \omega) = \bigotimes V / \langle u \otimes v - v \otimes u - \omega(u, v) \rangle$.

| | Even (V, ω) | odd ($\Pi V, g$) |
|--------------------|------------------------------------------------------------|-----------------------------------------------------------|
| graded alg. | $S(V^*)$ | ΛV^* |
| symplectic form | $\omega_{ij} d\xi^i \wedge d\xi^j$ | $g_{ij} d\xi^i \wedge d\xi^j$ |
| symmetries | $(S^2(V^*), \{\cdot, \cdot\}) \simeq \text{sp}(V, \omega)$ | $(\Lambda^2 V^*, \{\cdot, \cdot\}) \simeq \text{o}(V, g)$ |
| Moyal $*$ -product | $(S(V^*), *) \simeq \mathcal{W}(V^*, \omega^{-1})$ | $(\Lambda V^*, *) \simeq \mathbb{C}\text{I}(V^*, g^{-1})$ |
| Group action | $\text{Mp}(V, \omega) = \exp(\text{sp}(V, \omega))$ | $\text{Spin}(V, g) = \exp(\text{o}(V, g))$ |

| | Even (V, ω) | odd ($\Pi V, g$) |
|----------------|-------------------------------------------------------------------------------------------------------------------------------------|----------------------------------------------------------------------------------------------------------------------------------|
| Darboux coord. | $\omega = dp_i \wedge dx^i$ | $g = \varepsilon_i d\xi^i \wedge d\xi^i$; $\varepsilon_i = \pm 1$ Riemannian : $\varepsilon_i = 1$ |
| Lagrangian | $V = \langle x^i \rangle \oplus \langle p_i \rangle$ | $V \otimes \mathbb{C} = P \oplus \bar{P}$ $P = \left(\frac{1+iJ}{\sqrt{2}}\right) V$, $J^2 = -1$ Herm. str. |
| Polarized fct. | $S(\langle x^i \rangle)$ | ΛP^* |
| Quant. fct. | $S(\langle x^i \rangle)(1 + \langle p_i \rangle)$ | $\Lambda P^*(1 + \langle \bar{P}^* \rangle)$ |
| Rep. space | $\mathcal{H} = S(\langle x^i \rangle) \otimes \det^{\frac{1}{2}}$ $\text{End}(\mathcal{H}) \simeq \mathcal{W}(V^*, \omega^{-1})$ | $\mathcal{H} = \Lambda P^* \otimes \text{Ber}^{\frac{1}{2}}$ $\text{End}(\mathcal{H}) \simeq \mathbb{C}\text{I}(V^*, g^{-1})$ |

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| Darboux coord. | $\omega = dp_i \wedge dx^i$ | $g = \varepsilon_i d\xi^i \wedge d\xi^i \times \frac{\hbar}{2i}; \varepsilon_i = \pm 1$ Riemannian : $\varepsilon_i = 1$ |
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Symplectic supermanifold

Supermanifold := loc. ringed space $\mathcal{N} = (N, \mathcal{O})$ s.t. $\mathcal{O}|_U \simeq \mathcal{C}^\infty(\mathbb{R}^n) \otimes \Lambda \mathbb{R}^m$.

- Batchelor Theorem ('79) : $\mathcal{N} \longleftrightarrow (N, E)$, $\Pi E = (N, \Gamma(\Lambda E^*))$.
Example : $\mathcal{M} = T^*M \oplus \Pi TM$ **supercotangent** bundle over (M, g) .
- Rothstein Theorem ('91) : $(\mathcal{N}, \omega) \longleftrightarrow (N, \omega_0, E, g, \nabla)$.
 $\Rightarrow (\mathcal{M}, \omega)$ over (M, g) , given by $\omega = d\alpha$,

$$\color{red}\alpha = p_i dx^i + \frac{\hbar}{2i} g_{ij} \xi^i d^\nabla \xi^j.$$

Darboux coordinates (non-tensorial !) :

$$\tilde{\xi}^a = \theta_i^a \xi^i \quad \text{and} \quad \tilde{p}_i = p_i + \frac{\hbar}{2i} \omega_{iab} \tilde{\xi}^a \tilde{\xi}^b,$$

with $\theta_i^a \theta_j^b g^{ij} = \eta^{ab}$ and $\omega_{iab} = \partial_{\tilde{\xi}^a} \nabla_i \tilde{\xi}^b$.

Classical aspects of spin mechanics

The momentum of a rotation is :

$$\mathcal{J}_{x_{ij}} = p_i x_j - x_j p_i + \frac{\hbar}{i} \xi_i \xi_j,$$

hence the **spin** is $S_{ij} = \frac{\hbar}{i} \xi_i \xi_j$, generating $(\Omega^2_x(M), \{\cdot, \cdot\}) \simeq o(p, q)$.

Example

Rotating particle e.o.m.= Papapetrou equations = Hamiltonian flows of $g^{ij} p_i p_j$:

$$\begin{aligned}\dot{x}^j \nabla_j \dot{x}^i &= -\frac{1}{2} g^{ik} R(S)_{jk} \dot{x}^j, \\ \dot{x}^k \nabla_k S^{ij} &= 0.\end{aligned}$$

GQ of the supercotangent bundle

- Locally : merging of the GQ of $T^*\mathbb{R}^n$ and $\Pi\mathbb{R}^n$.
Explicit formulae in Darboux coord.
- Globally : we need a polarization in the fibers $\Pi T_x M$,
 \hookrightarrow almost Hermitian structure on (M, g) .
- Polarized functions : $\Lambda(T^{1,0}{}^* M)$, holomorphic diff. forms.
 $\Rightarrow \Gamma(\text{Cl}(M, g)) \simeq \text{End}(\Lambda(T^{1,0}{}^* M) \otimes \text{Ber}^{\frac{1}{2}})$, spinor bundle **S** !

Remark : coincide with usual construction of S over pseudo-Hermitian manifold :

- almost Hermitian structure \Rightarrow Spin^c-structure,
- + existence $\text{Ber}^{\frac{1}{2}} (= K^{\frac{1}{2}}) \Rightarrow$ Spin-structure,
- induced action of $u(n)$ on S is the natural one on
 $\Lambda(T^{1,0}{}^* M) \otimes \text{Ber}^{\frac{1}{2}}$.

$\text{Vect}(M)$ action on the supercotangent bundle

- Tensorial action : $\mathbb{L}_X = X^i \partial_i - p_j (\partial_i X^j) \partial_{p_i} + \xi^j (\partial_j X^i) \partial_{\xi^i}$.
- Hamiltonian action : $L_{\tilde{X}} \alpha = 0$, given by ?

Lemma

$$\tilde{X} = X^i \partial_i^\nabla + Y_{ij} \xi^j \partial_{\xi_i} - p_j \nabla_i X^j \partial_{p_i} + \frac{\hbar}{2i} \left(R_{lij}^k \xi_k \xi^l X^j - (\nabla_i Y_{kl}) \xi^k \xi^l \right) \partial_{p_i},$$

where Y is an arbitrary 2-form on M , depending linearly on X .

- $Y = 0 \Rightarrow \text{Vect}(M)$ -action iff ∇ is flat.

Lift by ∇ of Ham. action on T^*M , "moment map" : $J_X = p_i X^i$.

- $L_{\tilde{X}} \beta = f \beta$, where $\beta = g_{ij} \xi^i dx^j$, $d\beta$ odd sympl. form on ΠTM .

Then, conf(M, g)-action, $Y_{ij} = -\partial_{[i} X_{j]}$.

Moment map : $\mathcal{J}_X = p_i X^i + \frac{\hbar}{2i} \xi^j \xi^k \partial_{[k} X_{j]}$.

Vector fields action on the spinor bundle

We expect that an action of X on $\Gamma(S)$ is of the form

$$\nabla_X + Y_{ij}\gamma^i\gamma^j,$$

with Y as above. For the representation space $\Gamma(S)$, GQ is a Lie algebra morphism on : $Sb_{0,0} \oplus Sb_{1,0} \oplus Sb_{0,1} \oplus Sb_{0,2}$, where $Sb_{k,\kappa} = \text{Pol}_k(T^*M) \otimes \Omega_{\mathbb{C}}^\kappa(M)$. Hence,

Hamiltonian lift \iff Lie derivative of spinors .

Theorem

- $GQ(J_X) = \nabla_X$, *spinor covariant derivative* ;
- $GQ(\mathcal{J}_X) = \ell_X = \nabla_X + \frac{1}{4}\partial_{[k}X_{l]}\gamma^i\gamma^k$, *spinor Lie derivative (Kosmann'72)*.

Algebra of spinor differential operators D as a deformation of Sb

Spinless case : $\sigma_{p+q-1}([P, Q]) = \{\sigma_p(P), \sigma_q(Q)\}$.

Dirac operator : $\gamma^i \nabla_i$,

Lichnerowicz Laplacian : $[\gamma^i \nabla_i, \gamma^j \nabla_j] = \Delta + \frac{R}{4}$, second order diff. op.
⇒ usual filtration by order $D_0 \subset D_1 \subset \dots$ does not fit.

Idea : *Hamiltonian* filtration with ∇_i of order 1 and γ^i of order $\frac{1}{2}$.

$$[\nabla_i, \nabla_j] = R_{ijkl} \gamma^k \gamma^l \quad \text{and} \quad [\gamma^i, \gamma^j] = -2g^{ij}.$$

Theorem

D admits a filtration s.t. $D_K/D_{K-\frac{1}{2}} \simeq Sb_K = \bigoplus_{2k+\kappa=K} Sb_{k,\kappa}$ and
 $(Sb, \{\cdot, \cdot\})$ is the graded version of $(D, [\cdot, \cdot])$

$$\sigma_{K+L-1}([P, Q]) = \{\sigma_K(P), \sigma_L(Q)\}.$$

D as deformation of the $\text{conf}(M, g)$ -module Sb

Let us suppose that (M, g) is conformally flat, loc. $g_{ij} = F\eta_{ij}$.

- \mathbf{D} , adjoint action $\mathcal{L}_X A = [\ell_X, A]$, preserves Ham. and usual filtration.
- \mathbf{Sb} , Ham. action preserves Ham. graduation, on $Sb_{*, \kappa}$:

$$\tilde{X} = \mathbb{L}_X - \frac{\kappa}{n}(\nabla_i X^i) + \left(R_{lij}^k \xi_k \xi^l X^j - (\nabla_i Y_{kl}) \xi^k \xi^l \right) \partial_{p_i}.$$

As $\sigma(\text{GQ}(\mathcal{J}_X)) = \mathcal{J}_X$, we have $\sigma([L_x, \cdot]) = \{\mathcal{J}_X, \cdot\}$.
 $\Rightarrow (Sb, L)$ is the graded version of (D, \mathcal{L}) .

Proposition

$\mathcal{T} = \bigoplus_{\kappa} Sb_{*, \kappa} \otimes \mathcal{F}^{-\frac{\kappa}{n}}$ endowed with the tensorial action $\mathbb{L}^* = \mathbb{L}^{-\frac{\kappa}{n}}$ on $Sb_{*, \kappa}$ is the graded version of (Sb, L) and (D, \mathcal{L}) for the usual filtration.

Main result

Theorem

There exists a unique morphism of $\text{conf}(M, g)$ -modules preserving the principal symbol for the usual filtration

$$S_T^\delta : \mathcal{T}^\delta \rightarrow Sb^\delta,$$

*it is named the conformally equivariant **superization**.*

Theorem

There exists a unique morphism of $\text{conf}(M, g)$ -modules preserving the principal symbol for the Ham. filtration,

$$\mathcal{Q}^{\lambda, \mu} : Sb^\delta \rightarrow D^{\lambda, \mu},$$

*if $\mu - \lambda = \delta$. It is named the conformally equivariant **quantization**.*

Examples and applications

Examples :

① $S_T^0(J_X = p_i X^i) = \mathcal{J}_X = p_i X^i + \frac{\hbar}{2i} \xi^j \xi^k \partial_{[k} X_{l]}.$

② Dirac operator from De Rham differential : $\mathcal{Q} \circ S_T(p_i \xi^i) = \gamma^i \nabla_i.$

Examples and applications

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① $S_T^0(J_X = p_i X^i) = \mathcal{J}_X = p_i X^i + \frac{\hbar}{2i} \xi^j \xi^k \partial_{[k} X_{j]}.$

② Dirac operator from De Rham differential : $\mathcal{Q} \circ S_T(p_i \xi^i) = \gamma^i \nabla_i.$

Hamiltonian description of the Killing-Yano tensors, i.e. diff. κ -forms T s.t. $\nabla_{[i_0} T_{i_1 \dots i_\kappa]} = 0,$

$T \in \mathcal{T}_{0,\kappa}^{-\frac{1}{n}}$ s.t. $\{p_i \xi^i, S_T^0(\phi(T))\} = 0 \Leftrightarrow T$ is a **Killing-Yano** tensor.

Thanks !

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